# Amplitude analysis for exotic states at JPAC

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## Hadron Spectroscopy



Improvement needed! With great statistics comes great responsibility!

# Joint Physics Analysis Center

- Joint effort between theorists and experimentalists to work together to make the best use of the next generation of very precise data taken at JLab and in the world
- Created in 2013 by JLab & IU agreement
- It is engaged in education of further generations of hadron physics practitioners



## S-Matrix principles



 $A(s,t) = \sum_{l} A_{l}(s)P_{l}(z_{s})$  **Analyticity**  $A_{l}(s) = \lim_{\epsilon \to 0} A_{l}(s+i\epsilon)$ 

These are constraints the amplitudes have to satisfy, but do not fix the dynamics

Resonances (QCD states) are poles in the unphysical Riemann sheets

## Pentaquarks!





Quantum numbers  $J^{P} = \begin{pmatrix} 3^{-}, 5^{+} \\ \frac{3}{2}, \frac{5^{+}}{2} \end{pmatrix} \text{ or } \begin{pmatrix} 3^{+}, 5^{-} \\ \frac{3}{2}, \frac{5^{+}}{2} \end{pmatrix} \text{ or } \begin{pmatrix} 5^{+}, 3^{-} \\ \frac{5^{+}}{2}, \frac{3^{-}}{2} \end{pmatrix}$ Opposite parities needed for the interference to correctly describe angular distributions, low mass region contaminated by  $\Lambda^{*}$  (model dependence?)



## $P_c$ photoproduction

To exclude any rescattering mechanism, we propose to search the  $P_c(4450)$  state in photoproduction.



 $\langle \lambda_{\psi} \lambda_{p'} | T_r | \lambda_{\gamma} \lambda_p \rangle = \frac{\langle \lambda_{\psi} \lambda_{p'} | T_{\text{dec}} | \lambda_R \rangle}{M_r^2 - W^2 - \mathrm{i}\Gamma_r M_r} \frac{\langle \lambda_R | T_{\text{em}}^{\dagger} | \lambda_{\gamma} \lambda_p \rangle}{M_r^2 - W^2 - \mathrm{i}\Gamma_r M_r}$ 

### Hadronic part

- 3 independent helicity couplings,
  - $\rightarrow$  approx. equal,  $g_{\lambda_{\psi},\lambda_{p'}} \sim g$
- g extracted from total width and (unknown) branching ratio

Vector meson dominance relates the radiative width to the hadronic width

$$\Gamma_{\gamma} = 4\pi\alpha \, \Gamma_{\psi p} \left(\frac{f_{\psi}}{M_{\psi}}\right)^2 \left(\frac{\bar{p}_i}{\bar{p}_f}\right)^{2\ell+1} \times \frac{4}{6}$$

Hiller Blin, AP et al. (JPAC), PRD94, 034002

## **Background parameterization**

The background is described via an Effective Pomeron, whose parameters are fitted to high energy data from Hera



$$\lambda_{\psi}\lambda_{p'}|T_P|\lambda_{\gamma}\lambda_p\rangle = iA \left(\frac{s-s_t}{s_0}\right)^{\alpha(t)} e^{b_0(t-t_{\min})}\delta_{\lambda_p\lambda_{p'}}\delta_{\lambda_{\psi}\lambda_{\gamma}}$$

Asymptotic + Effective threshold

Helicity conservation

## Hiller Blin, AP et al. (JPAC), PRD94, 034002



## Pentaquark photoproduction



## Searching for resonances in $\eta\pi$

- The  $\eta\pi$  system is one of the golden modes for hunting hybrid mesons
- We build the partial waves amplitude according to the N/D method
- A. Jackura, et al. (JPAC & COMPASS), 1707.02848 see talk at 2:12pm



The denominator D(s) contains all the Final State Interactions constrained by unitarity  $\rightarrow$  universal The numerator n(s) depends on the exchanges  $\rightarrow$  process-dependent, smooth



## Searching for resonances in $\eta\pi$

The denominator D(s) contains all the FSI constrained by unitarity  $\rightarrow$  universal

natrix, ck for vanishing determinant

The numerator n(s) depends on the exchanges  $\rightarrow$  process-dependent, smooth

$$\rho_i(s)N_{ij}(s) = \frac{\lambda^{(2l+1)/2} \left(s, m_{\pi}^2, m_{\eta}^2\right)}{\left(s + \Lambda\right)^7}$$

## Searching for resonances in $\eta\pi$

The coupled channel analysis involving the  $\eta\pi$  and  $\eta'\pi$  for *P*- and *D*-wave is ongoing



# **Conclusions & prospects**

- We aim at developing new theoretical tools, to get insight on QCD using first principles of QFT (unitarity, analyticity, crossing symmetry, low and high energy constraints,...) to extract the physics out of the data
- Many other ongoing projects (both for meson and baryon spectroscopy, and for high energy observables), with a particular attention to producing complete reaction models for the golden channels in exotic meson searches



# BACKUP



## Production

- > 40 Research Papers (Phys.Rev., Phys.Lett, Eur.J. Phys.)
- ~120 Invited Talks and Seminars
- 0(10) ongoing analyses
- Summer Schools on Reaction Theory (IU, 2015 and 2017)
- Workshop "Future Directions in Hadron Spectroscopy" (JLab, 2014 and UNAM 2017)

| $\gamma N \to \pi \Delta$                                     | J. Nys <i>et al.,</i>        | arXiv:1710.09394             |
|---|------------------------------|------------------------------|
| FESR  | V. Mathieu <i>et al.,</i>    | arXiv:1708.07779             |
| $\pi N \to \eta \pi N$  | A. Jackura <i>et al.,</i>    | arXiv:1707.02848             |
| $\gamma N \rightarrow \eta N \text{ vs.} \rightarrow \eta' N$ | V. Mathieu <i>et al.,</i>    | arXiv:1704.07684             |
| <i>Z<sub>c</sub></i> (3900)                                   | A. Pilloni <i>et al.,</i>    | PLB772, 200                  |
| $\gamma N \rightarrow \eta N$                                 | J. Nys et al.,               | PRD95, 034014                |
| $\gamma p \rightarrow J/\psi p$                               | A. Blin <i>et al.,</i>       | PRD94, 034002                |
| $K N \rightarrow K N$   | C. Fernandez-Ramirez et al., | PRD93, 034029; PRD93, 074015 |
| $\gamma p \rightarrow \pi^0 p$                                | V. Mathieu <i>et al.,</i>    | PRD92, 074013                |
| $\pi N \to \pi N$   | V. Mathieu <i>et al.,</i>    | PRD92, 074004                |
| $\eta \rightarrow \pi^+  \pi^-  \pi^0$                        | P. Guo et al.,               | PRD92, 054016; PLB771, 497   |
| $\omega, \phi  ightarrow \pi^+ \pi^- \pi^0$                   | I. Danilkin <i>et al.,</i>   | PRD91, 094029                |
| $\gamma p \to K^+ K^- p$                                      | M. Shi <i>et al.,</i>        | PRD91, 034007                |
|   |                              |                              |

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## Interactive tools

- Completed projects are fully documented on interactive portals
- These include description on physics, conventions, formalism, etc.
- The web pages contain source codes with detailed explanation how to use them. Users can run codes online, change parameters, display results.

## http://www.indiana.edu/~jpac/

| Joint Physics Analysis Center |      |                 |                              |               |  |
|-------------------------------|------|-----------------|------------------------------|---------------|--|
|                               | HOME | PROJECTS        | PUBLICATIONS                 | LINKS         |  |
|                               |      | This project is | SFF National S<br>Foundation | Science<br>on |  |
|                               |      | $\pi N$         | $ ightarrow \pi N$           |               |  |
|                               |      |                 |                              |               |  |

### Formalism

The pion-nucleon scattering is a function of 2 variables. The first is the beam momentum in the laboratory frame  $p_{\rm lab}$  (in GeV) or the total energy squared  $s=W^2$  (in  ${\rm GeV^2}$ ). The second is the cosine of

#### Resources

- Publications: [Mat15a] and [Wor12a]
- SAID partial waves: compressed zip file
- C/C++: C/C++ file
- Input file: param.txt
   Output files: output0.txt , output1.txt , SigTot.txt , Observables0.txt , Observables1.txt
- Output files: output0.txt , output1.txt , SigTot.txt , Observables0.txt , Observables
   Contact person: Vincent Mathieu
- Contact person: Vincent r
   Last update: June 2016

The SAID partial waves are in the format provided online on the SAID webpage :

 $p_{
m lab} \quad \delta \quad \epsilon(\delta) \quad 1-\eta^2 \quad \epsilon(1-\eta^2) \quad {
m Re \, PW} \quad {
m Im \, PW} \quad SGT \quad SGR$ 

 $\delta$  and  $\eta$  are the phase-shift and the inelasticity.  $\epsilon(x)$  is the error on x. SGT is the total cross section and SGR is the total reaction cross section.

Format of the input and output files: [show/hide] Description of the C/C++ code: [show/hide]

### Simulation

| Range of the        | e running variab | le:   |      |      |   |
|---------------------|------------------|-------|------|------|---|
| $s$ in ${ m GeV}^2$ | (min max step)   | 1,2 ‡ | 6 ‡  | 0,01 | 1 |
| $p_{ m lab}$ in GeV | (min max step)   | 0,1 ‡ | 4 ‡  | 0,01 | ; |
| $\nu$ in GeV        | (min max step)   | 0,3 ‡ | 4 ‡  | 0,01 | 1 |
| $t~{ m in~GeV^2}$   | (min max step)   | -1 ‡  | 0 \$ | 0,01 | ; |

The fixed variable:

| in GeV <sup>2</sup> |  | 0  |
|---------------------|--|----|
| lab in GeV          |  | 5  |
| Start rese          |  | t) |

### Results



# Three-Body Unitarity

## Mai, Hu, Doring, AP, Szczepaniak, EPJA53, 9, 177

Original study by Amado/Aaron/Young

AAY(1968)

- 3-dimensional integral equation from unitarity constraint & BSE ansatz
- valid below break-up energies (E < 3m)
- analyticity constraints unclear

### One has to begin with asymptotic states



- *v* a general but cut-free (in the phys. region) function
- two-body interaction is parametrized by an "isobar"

= has definite QN and correct r.h.-singularities w.r.t invariant mass

• **S** and **T** are yet unknown functions

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M. Mai

## **Three-Body Unitarity**

A general ansatz for the Isobar-spectator interaction  $\rightarrow B \& \tau \text{ are unknown}!!!$ 







## **Three-Body Unitarity**

 $3 \rightarrow 3$  scattering amplitude is a 3-dimensional integral equation



- Imaginary parts (*B*,  $\tau$ , *S*) are fixed by **unitarity/matching** For simplicity  $v = \lambda$  (full relations available)

$$\tau(\sigma(k)) = (2\pi)\delta^{+}(k^{2} - m^{2})S(\sigma(k))$$
$$-\frac{1}{S(P^{2})} = \sigma(k) - M_{0}^{2} - \frac{1}{(2\pi)^{3}}\int d^{3}\ell \frac{\lambda^{2}}{2E_{\ell}(\sigma(k) - 4E_{\ell}^{2} + i\epsilon)}$$

$$\langle q|B(s)|p\rangle = -\frac{\lambda^2}{2\sqrt{m^2+\mathbf{Q}^2}\left(E_Q-\sqrt{m^2+\mathbf{Q}^2}+i\epsilon\right)}$$

- un-subtracted dispersion relation
- one- $\pi$  exchange in TOPT
- real contributions can be added to B



x



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## Triangle singularity



- Logarithmic branch points due to exchanges in the cross channels can simulate a resonant behavior, only in very special kinematical conditions (Coleman and Norton, Nuovo Cim. 38, 438)
- However, this effects cancels in Dalitz projections, no peaks (Schmid, Phys.Rev. 154, 1363)
- But the cancellation can be spread in different channels, you might still see peaks in other channels!